

A Neutron Diffraction Study of Ferroelectric, KFCT, $K_4Fe(CN)_6 \cdot 3D_2O$, above the Curie Temperature*

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A neutron diffraction study of deuterated potassium ferrocyanide trihydrate (KFCT) has been made above the ferroelectric Curie temperature. The monoclinic form, space group $A2/a$ (No. 15), has $a = 9.42$, $b = 17.01$, $c = 9.38$ Å, and $\beta = 90.03^\circ$ at 20°C . Since all of the large crystals of KFCT were twinned by a 90° rotation about the b axis, data from both overlapped and nonoverlapped reflections were used in the determination of the structure. Refinement by a modified least-squares method that accounted for the overlapped data led to a weighted R value on F^2 (3678 reflections) of 4.8%. The corresponding R on F is 4.9%. The disordered water molecules are of two types: one is disordered into two positions related by a twofold axis parallel to b : the other is disordered into two positions that are related by an approximate 90° rotation in the plane of the molecule. Both types of water molecules show $O-D \cdots O$ or $O-D \cdots N$ bonding. The $Fe(CN)_6^{4-}$ octahedra are regular and the potassium ions are loosely surrounded by O and N atoms. This analysis has elucidated the role of the water molecules in the ferroelectric switching mechanism. The [011] ferroelectric axis in this crystal can be explained on the basis of the dipole moment and orientation of the water molecule dipoles.

Introduction

An X-ray diffraction study of potassium ferrocyanide trihydrate (KFCT, hydrogen form) was made by Pospelov & Zhdanov (1947a). They reported that large single crystals were difficult to prepare because of twinning of the pseudotetragonal structure by a 90° rotation about b . Toyoda, Niizeki & Waku (1960) in a polarizing microscope study found at least three different types of twinning. However, all of the heavy atoms, except one water oxygen atom, were located by means of Patterson projections by Pospelov & Zhdanov (1947b). In an X-ray study Kiriyama, Kiriyama, Wada, Niizeki & Hirabayashi (1964) located the missing water oxygen atom.

Ferroelectric behavior in KFCT below -24.5°C (or -18°C for the deuterated form) was discovered by Waku, Masuro & Tanaka (1959) and by Waku, Hirabayashi, Toyoda & Iwasaki (1959), and further studied by Waku, Masuro, Tanaka & Iwasaki (1960). The nuclear magnetic resonance (n.m.r.) studies of KFCT were carried out by Blinc, Brenman & Waugh (1961), by Kiriyama, Kiriyama, Wada, Niizeki & Hirabayashi (1964) and by Tsang & O'Reilly (1965). These studies agreed that the onset of ferroelectric behavior is connected with an ordering of the hydrogen-bonded water molecules. Blinc, Brenman & Waugh (1961) considered that the electronic and ionic polarizabilities of K^+ and $Fe(CN)_6^{4-}$ as well as displacements of K^+ ions con-

tributed significantly to the internal field and spontaneous polarization. Structures for the ferroelectric phase were proposed in the three n.m.r. studies but these were all different.

Hazony, Earls & Lefkowitz (1968), based on Mössbauer studies, suggested that the mechanism is a mode instability (Cochran, 1960) and considered the K^+ ion to be vibrating against the $Fe(CN)_6^{4-}$ complex. They also stated that the water molecules have a substantial effect on the dynamics of the iron atom and provide the interaction between the balance of forces from K^+ to the $Fe(CN)_6^{4-}$ complex.

Since further clarification of the ferroelectric transition in KFCT depends, to a large extent, on the knowledge of the atomic positions, the present neutron diffraction study of the paraelectric phase was undertaken.

Experimental

The use of deuterated KFCT was preferred since (a) the deuterium atoms have a larger scattering length than hydrogen atoms for thermal neutrons, (b) the larger spin diffuse scattering by hydrogen results in considerable background, and (c) the measured linear absorption coefficient for the deuterated salt is $\mu = 0.44$ cm^{-1} , whereas it is 0.88 cm^{-1} for the hydrogen form.

The deuterated salt was prepared by dehydrating the normal salt at 110°C for 6 hours and dissolving the powder in 99.8% D_2O . Large crystals several mm on a side were grown by slow evaporation in a dry nitrogen atmosphere. The hydrogen content of the crystals was less than 0.5 mole per cent as checked by proton magnetic resonance.

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Most of the crystals had no marked extinction under polarized light; however, one crystal had good extinction throughout its volume, was optically biaxial, and appeared to be a monoclinic single crystal. The crystal was cooled to liquid nitrogen temperature to minimize extinction and to remove any tetragonal form (Toyoda, Niizeki & Waku, 1960). The crystal was shaped to a cylinder that had a diameter of 0.80 cm, a thickness of 0.37 cm, and the *b* axis along the cylinder axis.

Complete neutron diffraction data were collected at 20°C for 3678 reciprocal lattice points to $\sin \theta/\lambda = 0.82$, with $\lambda = 1.065 \text{ \AA}$. The data clearly indicated that the crystal was twinned (into two equal parts *A* and *B*) with *A* face centering and related by a 90° rotation about *b*. Since the monoclinic cell is practically tetragonal with $a \approx c$ and $\beta \approx 90^\circ$, the overlap of reciprocal lattice points was almost exact. Some points in the composite reciprocal lattice were superposed reflections, others were single reflections, while others had no intensity, corresponding to face-centered absences. (See Fig. 1.) The reflection types are defined in Table 1.

Table 1. Classification of reflection types in KFCT

Number of <i>hkl</i> 's	Type*	Observed as
1148	$\begin{cases} oee \\ eoo \end{cases}$	Single reflections (<i>A</i>)
2530		Superposed
3678	$\begin{cases} eee \\ ooo \end{cases}$	$I_{\text{obs}} = I_{hkl}(A) + I_{\bar{h}\bar{k}\bar{l}}(B)$
~ 1200	$\begin{cases} eeo \\ ooe \end{cases}$	Single reflections (<i>B</i>)
		Absent due to <i>A</i> face-centering
—	$\begin{cases} oeo \\ eoe \end{cases}$	

* *e* – even
o – odd

The measured intensities for 27 nonsuperposed reflections from part *A* and the corresponding reflections from part *B* were summed. Since the sums differed by less than 0.3%, the crystal appeared to be a 50:50 twin. The peak profiles were well-defined as shown in Fig. 2(a) with equal integrated intensities from the reflections of parts *A* and *B*. For comparison, the scans of a typical twinned crystal with a twin ratio quite different from 1:1 are shown in Fig. 2(b).

The cell dimensions of KFCT were determined by using the least-squares lattice constant program of Mueller, Heaton & Miller (1960). Results from three sets of neutron data from approximately 17 reflections are shown in Table 2. The single reflections serve to differentiate the *a* and *c* dimensions, whereas this is not true for the overlapped reflections. The *a* and *c* axes of Kiriyama, Kiriyama, Wada, Niizeki & Hirabayashi (1964) have been interchanged to agree with the conventional $a > c$ for monoclinic crystals as proposed by Donnay (1963).

A very small fragment of the same KFCT crystal used for neutron diffraction was used to collect limited X-ray diffraction intensities. The volume ratio *A*:*B* in this fragment was determined from X-ray intensities similar to that described above for neutrons. A value of approximately 3:1 was obtained as compared with 1:1 for the large crystal. This indicates that the twinning is not homogeneous.

Structure analysis

A three-dimensional Patterson synthesis of the neutron data gave the vector map of the superposed structures, and confirmed the positions for the $\text{Fe}(\text{CN})_6^{4-}$ proposed in the X-ray investigation by Pospelov & Zhdanov (1947b) and by Kiriyama, Kiriyama, Wada, Niizeki & Hirabayashi (1964). The positions of the K^+ ion and the water molecules were not revealed in the synthesis. The K^+ locations as determined by these two groups of investigators were assumed to be reliable. The water atoms were located in the present study by

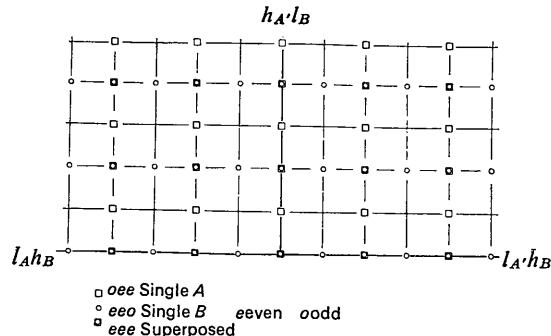


Fig. 1. Reciprocal lattice diagram for *k* even, illustrating the 90° twinning found in KFCT.

Table 2. Lattice constants of KFCT, space group *A*2/*a* (No. 15)

Reflections used	<i>a</i>	<i>b</i>	<i>c</i>	β
1	9.415 (7) Å	17.006 (13) Å	9.382 (7) Å	90.03 (2)°
2	9.420 (4)	17.040 (7)	9.396 (4)	90.08 (1)
3	9.428 (12)	16.937 (12)	9.437 (12)	90.07 (4)

- 1 Single reflections, part *A*.
- 2 Single reflections, part *B*.
- 3 Superposed reflections.

means of the following Fourier technique. The F^2 values were calculated for the K^+ and $Fe(CN)_6^{4-}$ contribution only; then the observed F^2 values for the overlapped reflections were divided up in the ratio of the twin components to derive approximations for $|F_{o(hkD)}|$ and $|F_{o(\bar{h}k\bar{D})}|$ as follows:

$$\begin{aligned} |F_{o(hkD)}| &= \sqrt{\frac{F_o^2 \times F_{c(hkD)}^2}{F_{c(hkD)}^2 + F_{c(\bar{h}k\bar{D})}^2}} \\ |F_{o(\bar{h}k\bar{D})}| &= \sqrt{\frac{F_o^2 \times F_{c(\bar{h}k\bar{D})}^2}{F_{c(hkD)}^2 + F_{c(\bar{h}k\bar{D})}^2}} \\ (F_o - F_c)_{hk\bar{D}} &= \frac{A_{hk\bar{D}}}{|A_{hk\bar{D}}|} \sqrt{\frac{F_o^2 \times F_{c(hkD)}^2}{F_{c(hkD)}^2 + F_{c(\bar{h}k\bar{D})}^2}} \\ &\quad - \frac{A_{hk\bar{D}}}{|A_{hk\bar{D}}|} \sqrt{F_{c(hkD)}^2} \\ (F_o - F_c)_{\bar{h}k\bar{D}} &= \frac{A_{\bar{h}k\bar{D}}}{|A_{\bar{h}k\bar{D}}|} \sqrt{\frac{F_o^2 \times F_{c(\bar{h}k\bar{D})}^2}{F_{c(hkD)}^2 + F_{c(\bar{h}k\bar{D})}^2}} \\ &\quad - \frac{A_{\bar{h}k\bar{D}}}{|A_{\bar{h}k\bar{D}}|} \sqrt{F_{c(\bar{h}k\bar{D})}^2}. \end{aligned}$$

The above Fourier coefficients were used to calculate Fourier and difference Fourier syntheses that clearly revealed the disordered water molecules. The water molecule I, near the twofold axis associated with O(1), was found to be disordered (including the oxygen atoms) into half-atoms on each side of the twofold axis (Fig. 3). The water molecule II, in the general location associated with O(2) O(3), was also disordered into two overlapped positions as shown in Fig. 4.

The initial model was refined on F^2 with the Kennicott (1963) modification of the least-squares program ORFLS, Busing, Martin & Levy (1962), which accounts for overlapped data. The number of F^2 values in the final cycles was 3678, and the value of

$$R = \frac{(\sum w(F_o - sF_c)^2)^{1/2}}{(\sum wF_o^2)^{1/2}}$$

was reduced to 4.6%; the value of w is given by the statistical weighting scheme of Evans (1961) where s is a scale factor.

The disordered positions for the oxygen atoms of water I were resolved in the Fourier syntheses and refined in the least-squares procedure; however, the half-oxygen atoms and superposed deuterium atoms of water II were too close to be resolved in the Fourier syntheses or refined by least-squares. Accordingly, the superposed oxygen and the superposed deuterium atoms of water II were refined as single atoms with double weight.

Fourier and difference Fourier syntheses were calculated from the final set of structure factors in the region of the water II molecule at $y=\frac{1}{2}$. The disordered atom sites in the overlapped peaks were estimated from the observed peak shapes. The theoretical shape for a half-oxygen or a half-deuterium atom was found by drawing the isotropic equivalent of a resolved half-deuterium [D(6) in this case]. The two theoretical half-atoms were then moved and their composite contours plotted until the composite curve fitted the observed peak shape. The observed and difference Fourier syntheses, on one-fifth the absolute scale, are shown at the $y=0.500$ level in Fig. 5(a) and (b) together with the positions proposed for the superposed half-atoms, O(2), O(3), D(3), and D(5). Bond lengths for the disordered molecules in water II were normal to within the estimated errors (Table 4). The separations O(2)-O(3) and D(3)-D(5) were 0.35 ± 0.25 and 0.65 ± 0.35 Å. In view of the disorder separations, it was not surprising that these overlapped atoms failed to refine as individual atoms.

An R value on F was derived by splitting the F_o^2 values of the overlapped reflections as described above

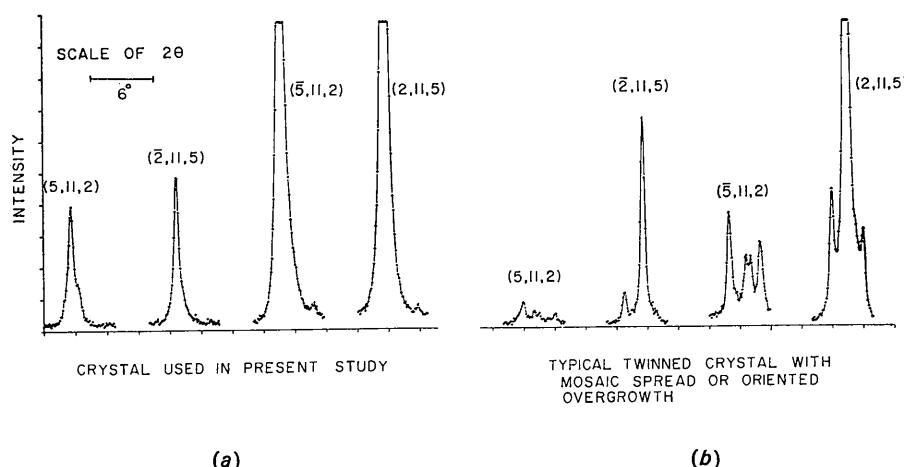


Fig. 2. Peak scans of 50:50 twin used in this study and scans of an unusable crystal exhibiting irregular peak shapes and a large mosaic spread.

The R value on F for the 3678 reflections was 4.9 %. The good agreement between observed and calculated structure factors justified our decision to proceed with a structural analysis of the twinned crystal. The final positional and thermal parameters are given in Table 3.

Extinction effects were not severe in the intensity data; only 17 single reflections and 38 doubled reflections were omitted from the least-squares refinement and are designated by an asterisk in Table 5.

The interatomic distances and angles are listed in Table 4. A list of the final observed and calculated structure factors is given in Table 5.

Discussion

Potassium ions and $\text{Fe}(\text{CN})_6^{4-}$ groups

Both potassium atoms K(1) and K(2) are loosely surrounded by oxygen and nitrogen atoms in an irregular fashion. K(1) has three oxygen atoms, one the disordered O(2)-O(3) atom, and four nitrogen atoms as nearest neighbors; K(2) is surrounded by one oxygen atom, the disordered O(2)-O(3) atom, and six nitrogen atoms. The K-O and K-N distances lie between 2.82 Å and 3.11 Å. (Table 4).

The octahedral $\text{Fe}(\text{CN})_6^{4-}$ group is nearly regular (see Table 4).

Water molecules

The water I and water II molecules are highly disordered, as shown in Figs. 3 and 4. Kiriyama, Kiriyama, Wada, Niizeki & Hirabayashi (1964) reported that dynamic disordering may occur with a reorienta-

tion frequency greater than 10^5 cps. The half-oxygen atoms of water I, disordered about the twofold axis, are separated by 0.86 Å; the half-deuterium atoms nearest the twofold axis are separated by 0.54 Å. The distances and angles in water I are normal (Table 4).

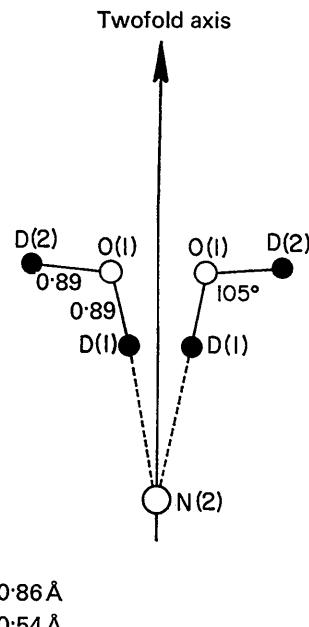


Fig. 3. Relative position of the disordered water molecule I found in KFCT.

Table 3. Positional and thermal parameters

The parameters and their standard deviations are multiplied by 10^4 . The temperature factor expression is

$$\text{T.F.} = \exp [-(\beta_{11} h^2 + 2\beta_{12} hk + 2\beta_{13} hl + \beta_{22} k^2 + 2\beta_{23} kl + \beta_{33} l^2)] .$$

Atom	x	y	z	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
Fe	2500	1765 (2)	0	0027 (2)	0016 (1)	0029 (2)	0	0001 (2)	0
N(1)	2500	-0053 (1)	0	0084 (4)	0012 (1)	0182 (6)	0	0047 (4)	0
N(2)	2500	3598 (2)	0	0067 (5)	0019 (5)	0099 (5)	0	0014 (4)	0
N(3)	1390 (2)	1748 (2)	6924 (2)	0060 (2)	0040 (1)	0038 (2)	0002 (2)	-0013 (1)	-0014 (1)
N(4)	4394 (2)	3219 (2)	6101 (2)	0044 (2)	0033 (1)	0071 (2)	-0013 (1)	0011 (1)	0006 (1)
C(1)	2500	0643 (2)	0	0062 (5)	0023 (1)	0066 (5)	0	0006 (5)	0
C(2)	2500	2915 (2)	0	0054 (4)	0012 (1)	0024 (4)	0	-0006 (4)	0
C(3)	1806 (2)	1774 (2)	8076 (2)	0039 (2)	0021 (1)	0034 (2)	0003 (1)	-0002 (1)	-0005 (1)
C(4)	4419 (2)	1778 (2)	9325 (2)	0033 (2)	0021 (1)	0040 (2)	-0008 (1)	0003 (1)	0002 (1)
K(1)	4350 (5)	1431 (3)	5963 (6)	0050 (5)	0034 (3)	0050 (5)	-0019 (2)	-0018 (3)	0001 (2)
K(2)	1576 (6)	3574 (3)	6907 (5)	0043 (5)	0049 (3)	0056 (6)	0007 (3)	-0015 (4)	-0021 (3)
O(1)	2276 (10)	0242 (4)	5391 (8)	0133 (10)	0016 (2)	0093 (8)	-0007 (4)	0023 (6)	0000 (3)
D(1)	2275 (11)	-0272 (4)	5176 (13)	0096 (12)	0030 (2)	0150 (12)	0016 (4)	0041 (8)	0037 (6)
D(2)	1702 (13)	0283 (6)	6144 (14)	0393 (26)	0037 (3)	0209 (16)	0012 (8)	0046 (16)	-0010 (6)
O(2)	0030 (130)	0000 (70)	7750 (130)	disordered atom	- coordinates from Fourier synthesis				
D(3)	0780 (190)	0060 (110)	8300 (190)	disordered atom	- coordinates from Fourier synthesis				
D(4)	-0784 (16)	-0214 (6)	8376 (12)	0554 (41)	0059 (4)	0168 (13)	-0096 (12)	-0055 (22)	0015 (7)
O(3)	0190 (130)	-0050 (70)	7500 (130)	disordered atom	- coordinates from Fourier synthesis				
D(5)	0690 (190)	-0230 (110)	8300 (190)	disordered atom	- coordinates from Fourier synthesis				
D(6)	0926 (17)	0281 (10)	7042 (25)	0391 (32)	0095 (8)	0524 (49)	-0104 (13)	-0010 (33)	0121 (16)
O(2)*	0164 (6)	-0013 (5)	7604 (7)	0191 (9)	0037 (2)	0257 (10)	-0024 (3)	-0100 (10)	0028 (3)
D(3)*	0675 (8)	-0087 (5)	8316 (5)	0337 (14)	0122 (5)	0122 (6)	0100 (7)	-0029 (9)	-0003 (4)

O(2)* is a composite of the superposed O(2) and O(3) in Water II.

D(3)* is a composite of the superposed D(3) and D(5) in Water II.

Water II, disordered by a 90° rotation in the plane of the molecule, is shown in Fig. 4. The half-oxygen atoms are disordered by 0.35 ± 0.25 Å and the half-deuterium atoms by 0.65 ± 0.35 Å. The dimensions of the overlapped molecules are normal, within the estimated errors (Table 4).

The arrangement of water molecules found by our neutron study is in agreement with the nuclear magnetic resonance (n.m.r.) study by Kiriya, Kiriya, Wada, Niizeki & Hirabayashi (1964) carried out at -170°C . The relation between the present results and the n.m.r. results is discussed in the section on ferroelectricity. The present results are in disagreement with the earlier n.m.r. study of Tsang & O'Reilly (1965) but agree with their later interpretation (O'Reilly & Tsang, 1967).

Deuterium bonding

(a) Water I

One deuterium, D(1), forms a bond with a length 2.84 Å with N(2) on the twofold axis, Fig. 3 and Table 4. The other deuterium, D(2), forms a bond with a length 2.83 Å with O(3) and has a much weaker possible interaction with N(3) with a distance of 3.05 Å. (See lower portion of Fig. 4).

(b) Water II

Deuterium bonding around water II is shown in Fig. 4. For the half-molecule D(3)-O(2)-D(4) the following possible O-D \cdots N distances are observed: 3.14, 3.19, and 3.26 Å. These distances suggest weak hydrogen bonding. In the half-molecule D(5)-O(3)-D(6), possible deuterium bonding distances of 2.83, 3.20, 3.25, and 3.31 Å occur.

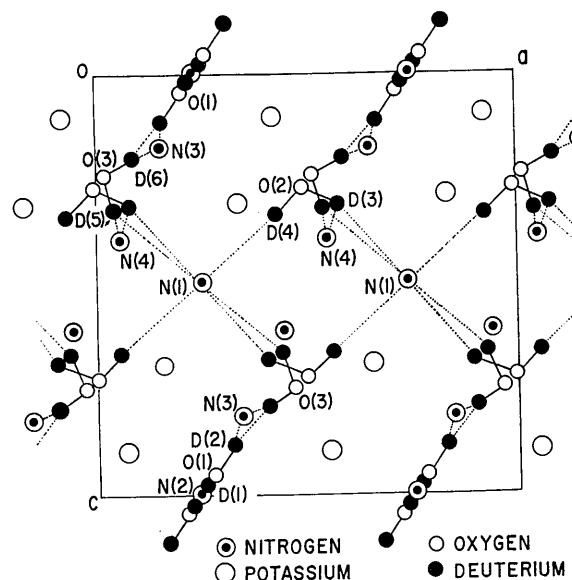


Fig. 4. The a - c projection of KFCT near $y=0.5$ showing the disordered water molecules and possible hydrogen bonding.

A complete distance and angle tabulation for possible deuterium bonding is given in Table 4. The deuterium bonding in the crystal is generally weak and the deuterium bonds are continually being made and broken in the disorder flipping, which is directly related to the ferroelectric switching.

Ferroelectricity in KFCT

If we combine the present neutron diffraction study of the paraelectric phase, which is disordered, with the n.m.r. study by Kiriya, Kiriya, Wada, Niizeki & Hirabayashi (1964) (carried out at -170°C , where the water molecules are rigid) the role of the water molecules in the ferroelectric transition becomes evident. The room temperature arrangement of disordered water molecules, Fig. 4, is equivalent to the

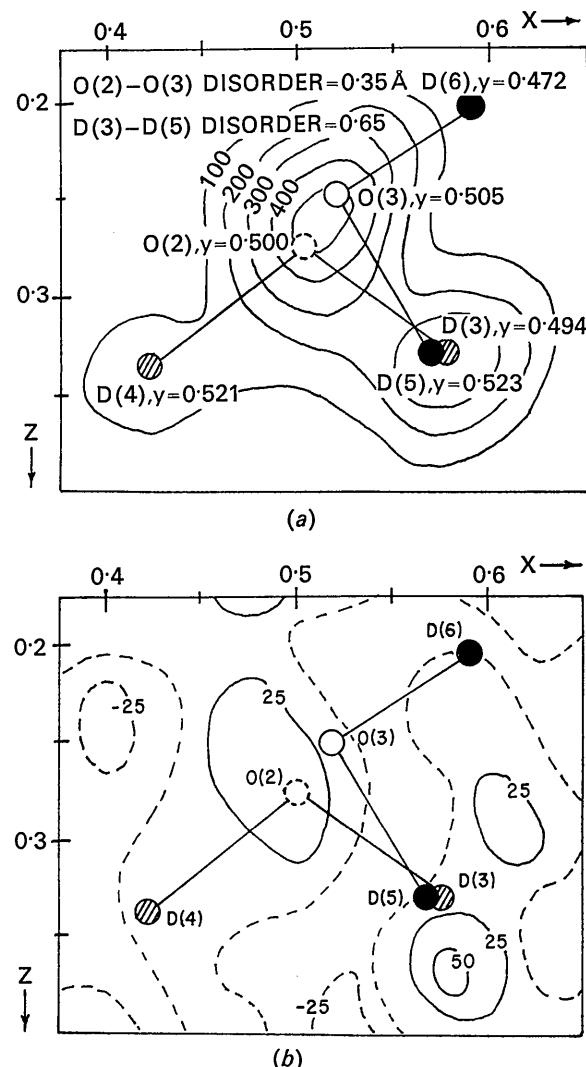


Fig. 5. (a) Observed Fourier section at $y=0.5$ showing disordered water molecule II. (b) A difference Fourier section at $y=0.5$ in the vicinity of water molecule II.

superposed ferroelectric structures proposed by Kiriyama, Kiriyama, Wada, Niizeki & Hirabayashi (1964). In one low-temperature ferroelectric state, the water molecule II splits into two non-equivalent water molecules associated with O(2) and O(3) and the space group becomes *Aa* (No. 9). One ferroelectric state is denoted by the solid arrow water dipoles corresponding to O(2) and O(3) in Fig. 6. The other ferroelectric state is denoted by the dashed arrows. The water molecule I is no longer associated with a twofold axis and assumes that one of the room temperature orientations which is dictated by steric requirements.

The definition of the ferroelectric axis arises naturally from the foregoing results since the dipole moment

of a water molecule bisects the D–O–D angle in the plane of the molecule. The ferroelectric axis in this crystal is the [101] since the moments associated with the water molecules of O(2) and O(3) are directed with one along [001] and the other along [100], as shown by the dashed arrows in Fig. 6. In a similar manner the solid arrows correspond to a [101] direction. The water I moment is forced to align in the same direction as the O(2), O(3) resultant. In the average structure above the Curie temperature, the water moments cancel.

Kiriyama, Kiriyama, Wada, Niizeki & Hirabayashi (1964) and Blinc, Brenman & Waugh (1961) agree that deviations from the ordered structure begin at -140°C .

Table 4. Interatomic distances and bond angles

Ferrocyanide groups			
Standard deviations ($\times 10^3$) for distances and ($\times 10^2$) for angles are in parentheses.			
Fe — C(1)	1.908 (4) Å	N(1)—C(1)—Fe	180.0°
Fe — C(2)	1.956 (4)	N(2)—C(2)—Fe	180.0
Fe — C(3)	1.920 (2)	N(3)—C(3)—Fe	177.38 (38)
Fe — C(4)	1.915 (2)	N(4)—C(4)—Fe	179.48 (55)
C(1)—N(1)	1.184 (3)	C(1)—Fe — C(3)	90.12 (3)
C(2)—N(2)	1.147 (3)	C(1)—Fe — C(4)	90.14 (3)
C(3)—N(3)	1.150 (2)	C(3)—Fe — C(4)	90.55 (10)
C(4)—N(4)	1.187 (2)	C(2)—Fe — C(3)	89.88 (3)
		C(2)—Fe — C(4)	89.86 (3)
Potassium contacts			
Standard deviations ($\times 10^3$) are in parentheses.			
K(1)—O(2)	*3.020 (100) Å	K(2)—O(2)	*2.880 (100) Å
K(1)—O(3)	*2.870 (100)	K(2)—O(3)	*3.060 (100)
K(1)—N(3)	2.979 (5)	K(2)—N(4)	2.824 (6)
K(1)—O(1)	2.862 (9)	K(2)—N(2)	3.029 (5)
K(1)—O(1)	2.836 (10)	K(2)—N(3)	3.050 (6)
K(1)—N(3)	2.848 (6)	K(2)—N(4)	3.026 (5)
K(1)—N(2)	3.101 (5)	K(2)—N(1)	3.067 (6)
K(1)—N(4)	3.043 (6)	K(2)—N(3)	3.109 (7)
* Superimposed half oxygen atoms.			
Water molecules			
Standard deviations ($\times 10^2$) for distance and ($\times 1$) for angles are in parentheses			
O(1)—D(1)	0.90 (1) Å	O(3)—D(5)	0.94 (10) Å
O(1)—D(2)	0.89 (1)	O(3)—D(6)	0.99 (10)
D(1)—O(1)—D(2)	105.00 (1)°	D(5)—O(3)—D(6)	101.00 (14)°
O(2)—D(3)	0.88 (10) Å		
O(2)—D(4)	1.03 (10)		
D(3)—O(2)—D(4)	108.00 (14)°		
Hydrogen bonds			
Standard deviations ($\times 10^2$) for distances and ($\times 1$) for angles are in parentheses			
O(1)—N(2)	2.84 (1) Å	O(2)—N(4)	3.26 (10) Å
D(1)—N(2)	1.96 (1)	D(3)—N(4)	3.19 (17)
O(1)—D(1)—N(2)	170.00 (1)°	O(2)—D(3)—N(4)	87.00 (12)°
O(1)—O(3)	2.83 (10) Å	O(3)—O(1)	2.83 (10) Å
D(2)—O(3)	1.99 (17)	D(6)—O(1)	2.00 (2)
O(1)—D(2)—O(3)	156.00 (14)°	O(3)—D(6)—O(1)	140.00 (6)°
O(1)—N(3)	3.05 (1) Å	O(3)—N(1)	3.20 (10) Å
D(2)—N(3)	2.61 (1)	D(5)—N(1)	2.35 (17)
O(1)—D(2)—N(3)	111.00 (1)°	O(3)—D(5)—N(1)	150.00 (14)°
O(2)—N(1)	3.14 (10) Å	O(3)—N(4)	3.25 (10) Å
D(3)—N(1)	2.28 (17)	D(5)—N(4)	2.70 (17)
O(2)—D(3)—N(1)	166.00 (15)°	O(3)—D(5)—N(4)	118.00 (14)°
O(2)—N(1)	3.19 (10) Å	O(3)—N(3)	3.31 (10) Å
D(4)—N(1)	2.29 (2)	D(6)—N(3)	2.54 (2)
O(2)—D(4)—N(1)	147.00 (14)°	O(3)—D(6)—N(3)	134.00 (6)°

Above -40°C a sudden increase in the reorientation frequency of the water molecules occurs. They also agree that no very large shifts in the heavier atoms occur at the Curie point of -18°C . Blinc, Brenman &

Waugh (1961) consider that electronic and ionic polarizabilities of K^+ and $\text{Fe}(\text{CN})_6^{4-}$ and some displacement of the K^+ ions could contribute significantly to the internal field and spontaneous polarization. The

Table 5. Observed and calculated structure factors ($\times 10$)

The first part of the Table contains the single reflections and the second part contains the overlapped reflections. The column headed $\text{F}^2\text{C}(\text{S}) = \text{F}^2\text{C}(\text{HKL}) + \text{F}^2\text{C}(\bar{\text{LKH}})$. Reflections indicated with an asterisk presumably suffer from extinction and were omitted from the final least-squares refinement.

	L	F^2O	F^2C																											
H^*	0	15	3	7	203	241	6	0	7	-9	42	29	-10	57	44	-3	889	681	12	40	39	-5	349	314	-7	145	75			
-9	0	14	6	9	643	500	10	11	17	-10	150	143	-11	117	115	-4	280	319	-5	149	165	-5	417	428	4	73	34			
-5	172	196	1	1	4	45	19	1	16	-5	637	537	-4	53	48	1	128	94	-10	1090	1137	9	45	45	-2	183	123			
-3	273	365	-5	0	39	8	38	87	1	420	515	7	71	79	5	91	92	-6	37	45	2	71	59	-2	228	216				
3	226	381	-1	0	39	4	294	389	5	1826	1840	1	21	15	7	1299	1083	-2	1856	1847	1	11	11	-7	356	395				
5	154	190	1	0	39	4	294	389	6	572	572	-6	62	50	1	153	141	-1	1090	1137	10	50	25	-2	222	225				
7	154	190	1	0	39	4	294	389	7	572	572	6	572	572	1	153	141	-1	1090	1137	9	57	70	-2	183	123				
9	15	14	7	7	46	17	2	670	904	11	113	113	1	153	141	1	128	94	1	10	308	105	1	11	11	-1	109	105		
11	15	14	7	7	46	17	2	670	904	12	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
13	15	14	7	7	46	17	2	670	904	13	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
15	15	14	7	7	46	17	2	670	904	14	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
17	15	14	7	7	46	17	2	670	904	15	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
19	15	14	7	7	46	17	2	670	904	16	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
21	15	14	7	7	46	17	2	670	904	17	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
23	15	14	7	7	46	17	2	670	904	18	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
25	15	14	7	7	46	17	2	670	904	19	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
27	15	14	7	7	46	17	2	670	904	20	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
29	15	14	7	7	46	17	2	670	904	21	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
31	15	14	7	7	46	17	2	670	904	22	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
33	15	14	7	7	46	17	2	670	904	23	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
35	15	14	7	7	46	17	2	670	904	24	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
37	15	14	7	7	46	17	2	670	904	25	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
39	15	14	7	7	46	17	2	670	904	26	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
41	15	14	7	7	46	17	2	670	904	27	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
43	15	14	7	7	46	17	2	670	904	28	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
45	15	14	7	7	46	17	2	670	904	29	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
47	15	14	7	7	46	17	2	670	904	30	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
49	15	14	7	7	46	17	2	670	904	31	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
51	15	14	7	7	46	17	2	670	904	32	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
53	15	14	7	7	46	17	2	670	904	33	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
55	15	14	7	7	46	17	2	670	904	34	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
57	15	14	7	7	46	17	2	670	904	35	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
59	15	14	7	7	46	17	2	670	904	36	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
61	15	14	7	7	46	17	2	670	904	37	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
63	15	14	7	7	46	17	2	670	904	38	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
65	15	14	7	7	46	17	2	670	904	39	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
67	15	14	7	7	46	17	2	670	904	40	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
69	15	14	7	7	46	17	2	670	904	41	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
71	15	14	7	7	46	17	2	670	904	42	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
73	15	14	7	7	46	17	2	670	904	43	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
75	15	14	7	7	46	17	2	670	904	44	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
77	15	14	7	7	46	17	2	670	904	45	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
79	15	14	7	7	46	17	2	670	904	46	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
81	15	14	7	7	46	17	2	670	904	47	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
83	15	14	7	7	46	17	2	670	904	48	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
85	15	14	7	7	46	17	2	670	904	49	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
87	15	14	7	7	46	17	2	670	904	50	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1	109	105
89	15	14	7	7	46	17	2	670	904	51	50	6	-11	1	0	6	8	1	128	94	1	10	308	105	1	11	11	-1		

Table 5 (cont.)

L	F ₀	F _C	F _{C(S)}	L	F ₀	F _C	F _{C(S)}	L	F ₀	F _C	F _{C(S)}	L	F ₀	F _C	F _{C(S)}	L	F ₀	F _C	F _{C(S)}	L	F ₀	F _C	F _{C(S)}							
H ^a	2 ^b	K ^c	1 ^d	2 ^e	10	291	106	269	H ^a	240	93	157	H ^a	5	52	8	5	11	H ^a	13	K ^c	158	H ^a	8	144	29	91			
	0 ^b	23	22	22	10	165	1	133		2	162	78	141		7	97	8	5	11		1	101	158		H ^a	1	378	217	400	
	2 ^b	175	140	181	H ^a	4	K ^c	14	H ^a	6	1078	450	1000	H ^a	9	1	46	124	H ^a	-1	K ^c	20	H ^a	0	151	120	165			
	0 ^b	917	287	909		0	125	586		4	118	26	34		1	142	7	7	149		0	127	143	189	H ^a	6	5	118	10	
	2 ^b	10	100	1	160		4	683	218	600	H ^a	6	2	16	26	H ^a	1	149	7	84		4	143	138	189	H ^a	0	58	58	112
N ^a	2 ^b	K ^c	18	18	10	18	5	5	H ^a	2	357	287	355	H ^a	2	201	64	221	H ^a	0	200	1233	2111	H ^a	6	85	7	26		
	4 ^b	0	0	37	H ^a	4	K ^c	18	H ^a	6	18	26	5	H ^a	2	201	63	158	H ^a	1	19	201	596	H ^a	2	181	55	54		
	4 ^b	156	18	93	H ^a	18	18	26		4	2	10	16		1	113	121	156	H ^a	2	0	18	37	H ^a	6	286	3	10		
	4 ^b	77	80	127	H ^a	4	K ^c	18	H ^a	6	51	2	4	H ^a	9	1	7	7	H ^a	-1	K ^c	20	H ^a	0	52	30	80			
	6 ^b	73	18	39	H ^a	4	K ^c	18	H ^a	6	51	2	4	H ^a	9	1	7	7	H ^a	0	103	0	5	H ^a	4	245	85	209		
H ^a	3 ^b	84	4	8	H ^a	6	138	137	201	H ^a	6	1	127	5	H ^a	9	1	7	7	H ^a	0	103	0	5	H ^a	4	283	193	258	
	2 ^b	303	147	164	H ^a	4	K ^c	18	H ^a	6	127	5	10	H ^a	9	1	7	7	H ^a	0	103	0	5	H ^a	4	283	193	258		
	6 ^b	373	121	209	H ^a	4	K ^c	18	H ^a	6	127	5	10	H ^a	9	1	7	7	H ^a	0	103	0	5	H ^a	4	283	193	258		
H ^a	2 ^b	K ^c	22	25	H ^a	6	102	18	96	H ^a	7	7725	3988	7827	H ^a	9	1	11	103	H ^a	-1	K ^c	24	H ^a	0	148	105	160		
	5 ^b	52	24	25	H ^a	6	145	37	133	H ^a	7	369	111	155	H ^a	9	1	148	105	H ^a	0	167	144	202	H ^a	4	101	244	452	
	2 ^b	211	70	89	H ^a	4	K ^c	20	H ^a	6	127	5	10	H ^a	9	1	148	105	H ^a	0	167	144	202	H ^a	4	101	244	452		
H ^a	3 ^b	K ^c	19	24	H ^a	6	20	50	H ^a	7	2857	1362	2873	H ^a	9	1	133	105	H ^a	0	167	144	202	H ^a	4	101	244	452		
	5 ^b	76	14	24	H ^a	6	223	152	342	H ^a	7	354	121	130	H ^a	9	1	128	113	H ^a	0	167	144	202	H ^a	4	101	244	452	
	5 ^b	310	102	105	H ^a	4	K ^c	20	H ^a	6	202	1	6	H ^a	9	1	148	105	H ^a	0	167	144	202	H ^a	4	101	244	452		
	7 ^b	101	33	107	H ^a	4	K ^c	20	H ^a	6	127	5	10	H ^a	9	1	148	105	H ^a	0	167	144	202	H ^a	4	101	244	452		
	7 ^b	1078	62	124	H ^a	4	K ^c	20	H ^a	6	127	5	10	H ^a	9	1	148	105	H ^a	0	167	144	202	H ^a	4	101	244	452		
	13 ^b	155	17	105	H ^a	4	K ^c	20	H ^a	6	127	5	10	H ^a	9	1	148	105	H ^a	0	167	144	202	H ^a	4	101	244	452		
H ^a	3 ^b	K ^c	21	21	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	15	40	H ^a	0	154	158	660	H ^a	4	154	158	660		
	5 ^b	245	10	211	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	5 ^b	40	12	12	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	5 ^b	70	8	21	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	5 ^b	9	169	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11 ^b	21	15	15	H ^a	5	K ^c	18	H ^a	6	154	158	660	H ^a	9	1	154	158	H ^a	0	154	158	660	H ^a	4	154	158	660		
	11																													

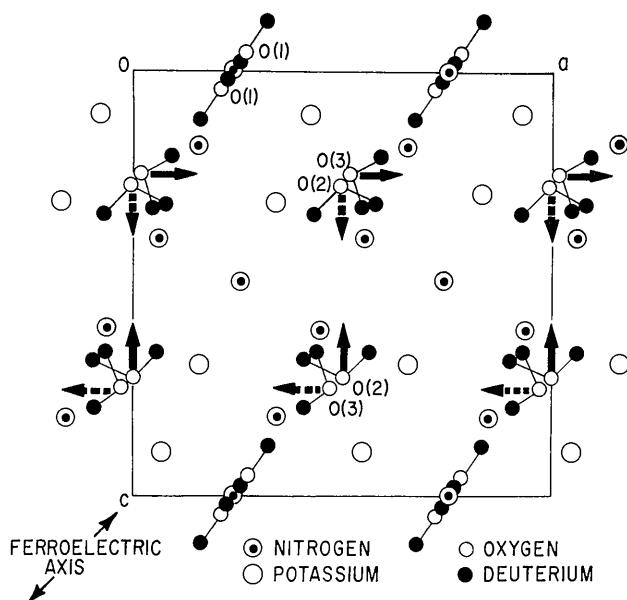


Fig. 6. KFCT water molecule layer at $y=0.5$ showing the ferroelectric axis. Water dipoles associated with one ferroelectric state are marked by the solid arrows, and with the other state by dotted arrows.

Further elucidation of the complex behavior of KFCT below the Curie temperature can best be obtained by a series of neutron diffraction studies at lower temperatures with an electrical field first along [101] and then a reversed field. This type of investigation was made by Bacon & Pease (1955) on KH_2PO_4 and would also allow a check on possible movements of the K^+ and the ferrocyanide ions.

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Short Communications

Contributions intended for publication under this heading should be expressly so marked; they should not exceed about 1000 words; they should be forwarded in the usual way to the appropriate Co-editor; they will be published as speedily as possible. Publication will be quicker if the contributions are without illustrations.

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Arcing and polytypism. By V. K. AGRAWAL, Department of Physics, Hastinapur College (University of Delhi), Moti Bagh, New Delhi-23, India

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An X-ray study of the arcing phenomenon in polytypic crystals of CdI_2 , PbI_2 and CdBr_2 has revealed that various polytypes are formed by the creation of stacking faults at different regular intervals during crystal growth.

An X-ray study of a large number of crystals of three polytypic substances, *viz.* 120 of CdI_2 (Agrawal & Trigunayat, 1969a,b), 60 of PbI_2 (Agrawal, Chadha & Trigu-

nayat, 1970a) and 12 of CdBr_2 (Agrawal & Trigunayat, 1970), has revealed that the arcing phenomenon, which consists of an extension of the diffraction spots into small